

Physics 618 2020

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Central Extensions ↙ Can have same G and non-iso-exts.

$$1 \rightarrow A \xrightarrow{\iota} \tilde{G} \xrightarrow{\pi} G \rightarrow 1$$

$$\iota(A) \subset Z(\tilde{G})$$

Thm: $\bar{E}(G, A)$ set of isom. classes of extensions of G by A

$$\bar{E}(G, A) \leftrightarrow \underline{H^2(G, A)} \quad \boxed{\begin{matrix} A \text{ is} \\ \alpha \\ \text{set} \end{matrix}}$$

$$H^2(G, A) = Z^2(G, A) / \sim$$

$$Z^2(G, A) : f : G \times G \rightarrow A$$

s.t.

$$f(g_1, g_2) f(g_1 g_2, g_3) = f(g_2, g_3) f(g_1, g_2 g_3)$$

$$\sim : \hat{f}(g_1, g_2) = f(g_1, g_2) \frac{t(g_1) t(g_2)}{t(g_1 g_2)}$$

for some $t : G \rightarrow A$.

But $H^2(G, A)$ has a $\neq$ group structure.

$\overline{E}(G, A)$ can be intrinsically defined.

$$\mathcal{E}_1: 1 \rightarrow A \rightarrow \tilde{G}_1 \rightarrow G \rightarrow 1$$

$$\mathcal{E}_2: 1 \rightarrow A \rightarrow \tilde{G}_2 \rightarrow G \rightarrow 1$$

$$1 \rightarrow A \rightarrow \tilde{G}_{12} \rightarrow G \rightarrow 1$$

$$\mu(a_1, a_2) = a_1 a_2$$

$$\uparrow \mu$$

$$1 \rightarrow A \times A \rightarrow \hat{G}_{12} \rightarrow G \rightarrow 1$$

$$\mathcal{E}_1 \times \mathcal{E}_1: 1 \rightarrow A \times A \xrightarrow{2, \times 2} \tilde{G}_1 \times \tilde{G}_2 \xrightarrow{\pi_1 \times \pi_2} G \times G \rightarrow 1$$

$$\mathcal{E}_1 \cdot \mathcal{E}_2 = \mu_* \Delta^* (\mathcal{E}_1 \times \mathcal{E}_2)$$

$\mathcal{E}_1 \cdot \mathcal{E}_2$ should be an extension of G by A

Example 2 Extensions of $\mathbb{Z}_p$ by $\mathbb{Z}_p$

$$1 \rightarrow \mathbb{Z}_p \rightarrow \tilde{G} \rightarrow \mathbb{Z}_p \rightarrow 1$$

$\uparrow$ order = p^2

$$\Rightarrow \tilde{G} = \mathbb{Z}_p \times \mathbb{Z}_p \quad \left(\begin{array}{l} \text{in this} \\ \text{case the} \\ \text{extension is} \\ \text{trivial} \end{array} \right)$$

or $\mathbb{Z}_{p^2}$: There are several inequivalent extensions.

But : $H^2(\mathbb{Z}_p, \mathbb{Z}_p) \cong \mathbb{Z}_p$

$$1 \rightarrow \mathbb{Z}_p \xrightarrow{\iota} \mathbb{Z}_{p^2} \xrightarrow{\pi} \mathbb{Z}_p \rightarrow 1$$

$\nwarrow \quad \nearrow$
Conceptually different

So: Introduce notation:

$$\mathbb{Z}_p = \langle \sigma_1 \mid \sigma_1^p = 1 \rangle$$

$$\mathbb{Z}_{p^2} = \langle \alpha \mid \alpha^{p^2} = 1 \rangle$$

$$\mathbb{Z}_p = \langle \sigma_2 \mid \sigma_2^p = 1 \rangle$$

π is surjective

$$\begin{array}{ccccccc} \hookrightarrow \mathbb{Z}_p & \xrightarrow{2} & \mathbb{Z}_{p^2} & \xrightarrow{\pi} & \mathbb{Z}_p & \longrightarrow & 1 \\ \langle \sigma_1 \rangle & & \langle \alpha \rangle & & \langle \sigma_2 \rangle & & \end{array}$$

$$2(\sigma_1) = \alpha^x$$

$$2(1) = 2(\sigma_1^p) = 1 \implies \alpha^{p^x} = 1$$

$$\implies x = 0 \pmod{p}$$

2 injective so
must be of form:

$$\boxed{2_k(\sigma_1) = \alpha^{kp}}$$

where $k \in \mathbb{Z}_p^*$, $1 \leq k \leq p-1$

$$(k, p) = 1.$$

Similar: π must be of the form

$$\pi_r(\alpha) = \sigma_2^r \quad 1 \leq r \leq p-1$$
$$r \in \mathbb{Z}_p^*$$

$$\ker \pi_r \quad \pi_r(\alpha^l) = \sigma_2^{lr} = 1$$

$$r \in \mathbb{Z}_p^* \quad lr = 0 \pmod{p}$$

$$\therefore l = 0 \pmod{p}$$

$$\ker \pi_r = \{1, \alpha^p, \alpha^{2p}, \dots, \alpha^{(p-1)p}\}$$

$$\underline{\text{Check}} = \text{im } \mathcal{Z}_k$$

Conc: We've found the most general possibilities for $\mathcal{Z}, \pi$.

$$\mathcal{Z} = \mathcal{Z}_k \quad \text{for some } k \in \mathbb{Z}_p^*$$

$$\pi = \pi_r \quad \text{for some } r \in \mathbb{Z}_p^*$$

Compute a cocycle —

Choose a section:

$$1 \rightarrow \mathbb{Z}_p \xrightarrow{2_k} \mathbb{Z}_{p^2} \xrightarrow{\pi_r} \mathbb{Z}_p \rightarrow 1$$

$\downarrow \text{---} \downarrow$
 $S \qquad \qquad \qquad \langle \sigma_2 \rangle$

$$S(1) = 1$$

$$S(\sigma_2) = \alpha^x$$

$$\pi_r \circ S(\sigma_2) = \sigma_2$$

$$\therefore \pi_r(\alpha^x) = \sigma_2^{rx} = \sigma_2$$

$$\therefore rx = 1 \pmod{p}$$

Choose simplest solution $r \in \mathbb{Z}_p^*$

$$\exists! \quad 1 \leq r^* \leq p-1 \quad 1 \leq r \leq p-1$$

$$\text{s.t.} \quad rr^* = 1 \pmod{p}.$$

$$S(\sigma_2) = \alpha^{r^*}$$

$$S(\sigma_2^2) = \alpha^{2r^*}$$

$\vdots$

$$S(\sigma_2^{p-1}) = \alpha^{(p-1)r^*}$$

$$\begin{aligned}
 S(1) &= 1 \\
 S(\sigma_2) &= \alpha^{r^*} \\
 S(\sigma_2^2) &= \alpha^{2r^*} \\
 &\vdots \\
 S(\sigma_2^{p-1}) &= \alpha^{(p-1)r^*}
 \end{aligned}$$

$$S(\sigma_2^*) \in \mathbb{Z}_{p^2} \langle \alpha \rangle$$

Now $S(\sigma_2^p)$?

$$\begin{aligned}
 \text{group homom. requires } S(\sigma_2^{p-1})S(\sigma_2) \\
 = \alpha^{(p-1)r^*} \alpha^{r^*} = \alpha^{pr^*}
 \end{aligned}$$

$$\text{but } \sigma_2^p = 1 \quad S(1) = 1.$$

We're stuck: we can't choose S to be a group homomorphism.

Extension is non-split

Calculate the cocycle:

$$S(\sigma_2^x) S(\sigma_2^y) S(\sigma_2^{x+y})^{-1}$$

$$= \begin{cases} 1 & x+y \leq p-1 \\ \alpha^{r^*p} & x+y \geq p \end{cases}$$

$\therefore$

$$z_k(\sigma_1) = \alpha^k \quad k \in \mathbb{Z}_p^*$$

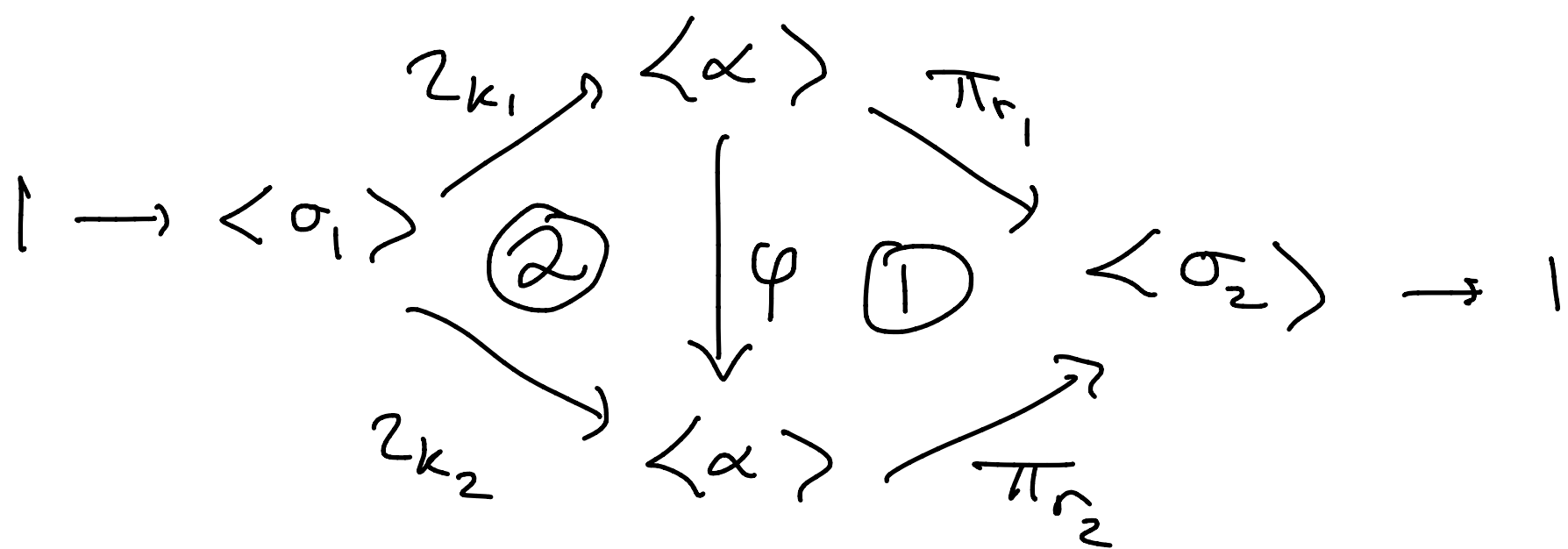
$$f_{k,r}(\sigma_2^x, \sigma_2^y) = \begin{cases} 1 & x+y \leq p-1 \\ \sigma_1^{k^*r^*} & x+y \geq p \end{cases}$$

$$A = \mathbb{Z}_p = \langle \sigma_1 \rangle$$

where $k^*k = 1 \pmod{p}$

→ This is a cocycle.

Notes: Show how trying to trivialize with a coboundary won't work.



isom. of extensions is an iso of groups making the diagram commute.

$$\varphi(\alpha) = \alpha^y \quad y \text{ rel prime to } p$$

① commutes

$$\pi_{r_2} \circ \varphi(\alpha) = \pi_{r_1}(\alpha)$$

$$\Rightarrow r_2 y = r_1 \pmod{p}$$

② commutes: $\varphi \tau_{k_1}(\sigma_1) = \tau_{k_2}(\sigma_2)$

$$\Rightarrow k_1 p y = k_2 p \pmod{p^2}$$

$$\Leftrightarrow k_1 y = k_2 \pmod{p}.$$

$$r_2 y = r_1 \pmod{p}$$

$$k_1 y = k_2 \pmod{p}$$

$$\text{and } y \in \mathbb{Z}_p^*$$

$$\Rightarrow k_1 r_1 = k_2 r_2 \pmod{p}$$

$[f_{k,r}]$ is determined by $k \cdot r$ or equiv. $k^* r^*$

$$[f_{k_1, r_1}] [f_{k_2, r_2}] \Leftrightarrow \underbrace{(k_1^* r_1^*) + (k_2^* r_2^*)}_{\pmod{p}}$$

$\mathbb{Z}_p$ ring $\mathbb{Z}_p^* \subset \mathbb{Z}_p$ group of units.

$$H^2(\mathbb{Z}_p, \mathbb{Z}_p) \cong \mathbb{Z}_p$$

$1 \leq (kr)^* \leq p-1$ $p-1$ extensions
with $\tilde{G} = \mathbb{Z}_{p^2}$ $\tilde{G} = \mathbb{Z}_p \times \mathbb{Z}_p / [f] = 1$.

Example 4 p prime $G \ni \vec{x}$

$$1 \rightarrow \mathbb{Z}_p \rightarrow \tilde{G} \rightarrow \underbrace{\mathbb{Z}_p \oplus \dots \oplus \mathbb{Z}_p}_{k \text{ summands}} \rightarrow 1$$

Write additively

Cocycle condition

$$f: G \times G \rightarrow \mathbb{Z}/p\mathbb{Z}$$

$$\begin{aligned} f(\vec{x}, \vec{y}) + f(\vec{x} + \vec{y}, \vec{z}) \\ = f(\vec{x}, \vec{y} + \vec{z}) + f(\vec{y}, \vec{z}) \end{aligned}$$

Satisfy that with

$$f(\vec{x}, \vec{y}) = \underline{\underline{A_{ij} x_i y_j}}$$

Just plug it in

$$A \in \text{Mat}_{k \times k}(\mathbb{Z}_p)$$

Change by a coboundary

$$f(\vec{x}, \vec{y}) \rightarrow f(\vec{x}, \vec{y}) + q(\vec{x} + \vec{y}) - q(\vec{x}) - q(\vec{y})$$

$q \sim$ coboundary. $q: G = (\mathbb{Z}/p\mathbb{Z})^k$

$$q(\vec{x}) = q_{ij} x_i x_j \rightarrow A = \mathbb{Z}/p\mathbb{Z}$$

$$A_{ij} \rightarrow A_{ij} - (q_{ij} + q_{ji})$$

$P=2$ put $A_{ij} = 0 \quad i > j \leftarrow$

$$A_{ji} = 0 \text{ or } 1$$

gauge inv.

$$A_{ii} \rightarrow A_{ii} - \underbrace{(2q_{ii})}_0 = \underline{\underline{A_{ii}}}$$

$$\frac{1}{2} k(k+1) \text{ indep cocycles.}$$

$$\mathbb{Z}_2^{\oplus \frac{1}{2} k(k+1)}$$

If p is an odd prime.

$$A_{ij} \rightarrow A_{ij} - 2(q_{ij} + q_{ji})$$

So $(\mathbb{Z}/p\mathbb{Z})^{\frac{1}{2}k(k-1)}$ diml
space
of cohomology
classes.

Using topology
one computes
(Künneth theorem)

$$H^2(\mathbb{Z}_p^{\oplus k}, \mathbb{Z}_p) \cong \mathbb{Z}_p^{\frac{1}{2}k(k+1)}$$

But we just constructed p

nontrivial cocycles $H^2(\mathbb{Z}_p, \mathbb{Z}_p)$

not of the form $A_{ij}x_iy_j$

Example 5:

$$1 \rightarrow \mathbb{Z}_2 \rightarrow \tilde{G} \rightarrow \mathbb{Z}_2 \oplus \mathbb{Z}_2 \rightarrow 1$$

$$H^2(\mathbb{Z}_2 \oplus \mathbb{Z}_2, \mathbb{Z}_2) \cong \mathbb{Z}_2 \oplus \mathbb{Z}_2 \oplus \mathbb{Z}_2$$

Distinct groups in the middle: 4 types

$$1 \rightarrow \mathbb{Z}_2 \rightarrow \mathbb{Z}_2 \times \mathbb{Z}_2 \times \mathbb{Z}_2 \rightarrow \mathbb{Z}_2 \times \mathbb{Z}_2 \rightarrow 1$$

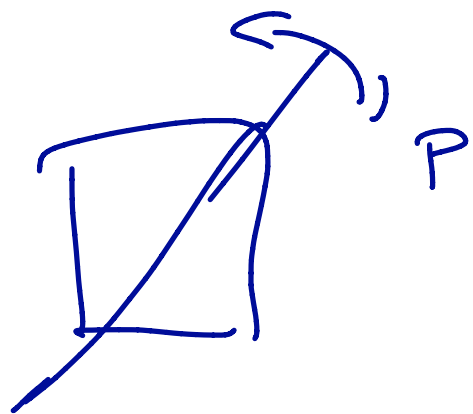
$$1 \rightarrow \mathbb{Z}_2 \rightarrow \mathbb{Z}_2 \oplus \mathbb{Z}_4 \rightarrow \mathbb{Z}_2 \oplus \mathbb{Z}_2 \rightarrow 1$$

$$1 \rightarrow \mathbb{Z}_2 \rightarrow Q \rightarrow \mathbb{Z}_2 \oplus \mathbb{Z}_2 \rightarrow 1$$

$$1 \rightarrow \mathbb{Z}_2 \rightarrow D_4 \rightarrow \mathbb{Z}_2 \oplus \mathbb{Z}_2 \rightarrow 1$$

$$Q = \{ \pm 1, \pm i\sigma^1, \pm i\sigma^2, \pm i\sigma^3 \}$$

$$D_4 = \{ 1, R, R^2, R^3, P, PR, PR^2, PR^3 \}$$



$$R = R(\pi/2)$$

Make an
Explicit choice of section $\Rightarrow$
explicit bilinear form $A_{ij} x_i y_j$.

8 extensions

4 isom. of groups $\hat{G}$.

Exercises:

1. Trivializable central extension
is a direct product $\tilde{G} = A \times G$

Not true of split (noncentral)
extensions.

2. D_4 vs. Q .

3. Group Commutator Criterion.

$$1 \rightarrow A \rightarrow \hat{G} \xrightarrow{\pi} G \rightarrow 1$$

$\tilde{g}_1, \tilde{g}_2 \xleftarrow{g_1 g_2}$

If $g_1, g_2 \in G$ are a commuting
pair $g_1 g_2 = g_2 g_1$ and there are

lifts $\tilde{g}_1, \tilde{g}_2$ s.t.

$$\pi(\tilde{g}_1) = g_1 \quad \pi(\tilde{g}_2) = g_2$$

if $\underline{[\tilde{g}_1, \tilde{g}_2]} = z(a) \quad a \neq 1$

Then the extension is nontrivial.



$$H^2(\text{Aut}(\text{QM}), U(1)) = ?$$

$$H^2(O(3), U(1))$$

which coho?

Single Qbit

Continuous.

$$H^*(BO(3), \mathbb{Z})$$

known

$$\langle w_1, w_2, \dots, p_1, \dots \rangle$$

$$w_1^2, w_2^2$$

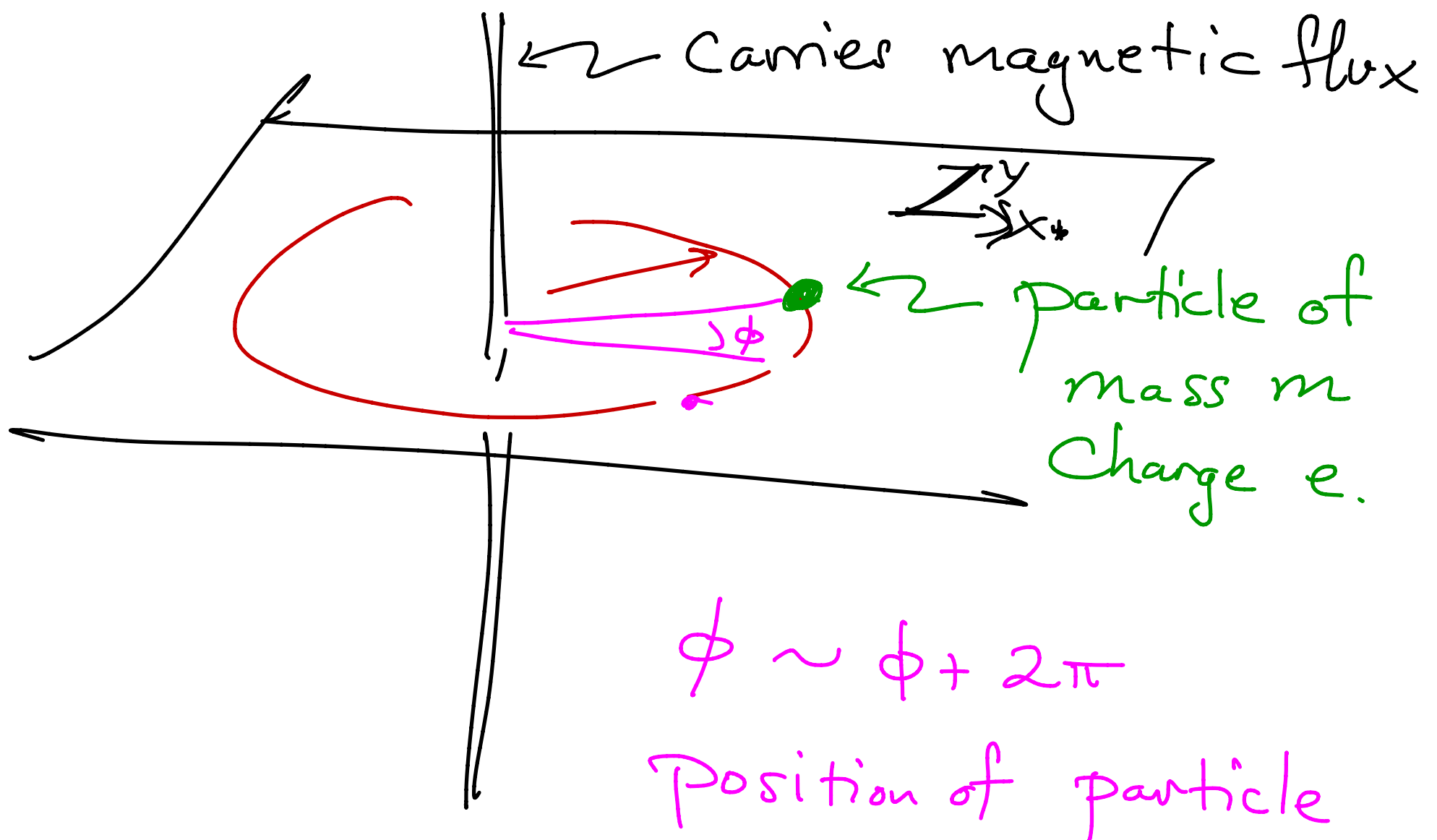
2-torsion

Pure states $S^2 = \partial$ Bloch Sphere

F-S metric: round metric.

$$\text{Aut}(\text{QM}) = \text{Isom}(S^2, \text{round metric}) = O(3).$$

Quantum Mechanics Of A Charged Particle On a Ring Surrounding a Solenoid:



$$S = \int \frac{1}{2} m r^2 \dot{\phi}^2 dt$$

$$= \int \frac{1}{2} I \dot{\phi}^2 dt \quad \text{w/out charge}$$

$$S = \int \frac{1}{2} I \dot{\phi}^2 + \int e A \quad \begin{array}{l} \swarrow \text{gauge potential} \\ \searrow \text{worldline} \end{array}$$

As in Aharonov-Bohm.

In general $x^\mu(t)$ moving through a space with gauge field $A_\mu(x)$ particle has charge e mass m

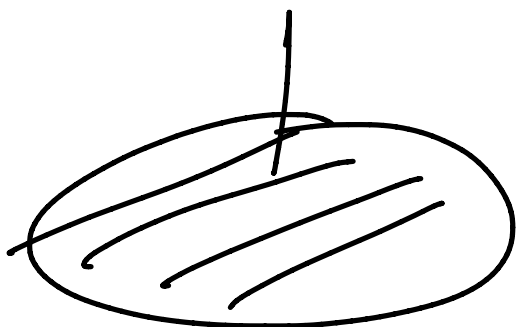
$$S' = \int \frac{1}{2} m \overset{\substack{\uparrow \\ \text{metric on space}}}{g_{\mu\nu}(x(t))} \dot{x}^\mu(t) \dot{x}^\nu(t) dt + \int \underbrace{e A_\mu(x(t)) \dot{x}^\mu(t)}_{\substack{\uparrow \\ \text{A-B phase}}} dt$$

$$\int dx(t) e^{i S}$$

$$S = \int \frac{1}{2} I \dot{\phi}^2 + \int \underbrace{e A_\mu(x(t)) \dot{x}^\mu(t)}_{\text{A-B phase}} dt$$

A 1-form

$$F = dA$$



$$\int F = B$$

1

$$\Rightarrow A = \frac{B}{2\pi} d\phi$$

$$A_z = 0 \quad A_r = 0 \quad A_\phi = \frac{B}{2\pi}$$

$$S = \int \frac{1}{2} I \dot{\phi}^2 dt + \frac{eB}{2\pi} \phi dt$$

Classical Symmetries

Quantum Symmetries

limit $\phi \rightarrow -\phi$
not

Equations of motion: $\ddot{\phi} = 0$.

N.B. ! B-term did not contribute.

$$\int \frac{eB}{2\pi} \dot{\phi} dt$$

total derivative

example of
"θ-term"
"topological term"

does not affect classical physics

does affect quantum physics

Classical eq's of motion $\ddot{\phi} = 0$

Symmetries $\phi \longrightarrow \phi + \alpha \quad \alpha \sim \alpha + 2\pi$

$$R(\alpha) : e^{i\phi(t)} \longrightarrow e^{i\phi(t)} e^{i\alpha}$$

$$= \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix}$$

$$P : \phi \longrightarrow -\phi \quad \text{reflection in } xy \text{ plane}$$

Symmetry of eqs of motion by not of the top. term.

$$R(\alpha) R(\beta) = R(\alpha + \beta) \quad SO(2)$$

$$P^2 = 1$$

$$P R(\alpha) P = R(-\alpha)$$

$$\boxed{O(2)}$$

Classical Symmetry
Group

Compute Hamiltonian:

Canonical momentum conj. to ϕ

$$L = \frac{\delta S}{\delta \dot{\phi}} = I \dot{\phi} + \frac{eB}{2\pi}$$

$$\int H dt = \int L \dot{\phi} dt - S'$$

$$= \int \frac{1}{2I} \left(L - \frac{eB}{2\pi} \right)^2$$

quant: $L = -i\hbar \frac{\partial}{\partial \phi}$ acting

on L^2 -wavefunctions of ϕ

$$L = -i\hbar \frac{\partial}{\partial \phi} + \underbrace{\text{const.}}_{=0}$$

$$H_B = \frac{\hbar^2}{2I} \left(-i \frac{\partial}{\partial \phi} - B \right)^2$$

$$B = \frac{eB}{2\pi\hbar}$$

Diagonalize $\mathcal{H} = L^2 (S^1)$

ON
Complete set: $\psi_m = \frac{1}{\sqrt{2\pi}} e^{im\phi} \quad m \in \mathbb{Z}$

also eigenvectors of H

$$H_B \psi_m = E_m \psi_m$$

$$E_m = \frac{\hbar^2}{2I} (m - B)^2$$

1-parameter
family of
Hamiltonians

$$\boxed{B \rightarrow B + c}$$

Remarks: $B^{\text{eff}} = B + c$
 $\boxed{c=0}$

does not change
the family.

1. Topological term matters because
the observable spectrum is
shifted by $B = eB/2\pi\hbar$

2. Important that $m \in \mathbb{Z}$
because $\phi \sim \phi + 2\pi$

Action $S[\phi(t)]$ also makes

Sense for $\phi(t) \in \mathbb{R}$

without the identification $\phi \sim \phi + 2\pi$

But then m would not be quantized and we could shift away $\mathcal{B}$ (or $\mathcal{B}^{\text{eff}}$).

Topological term would have no effect in the quantum theory as well.

Need to have some topology.

$$\text{Tr} \left[\frac{1}{H_{\mathcal{B}} - z} \right] = R(z)$$

very useful. $\frac{1}{H_{\mathcal{B}} - z}$ Resolvent

Next time:

After some further remarks
We'll look at the quantum symmetries.